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Optical and Radiative Properties of Ice Clouds

Lei Bi & Ping Yang

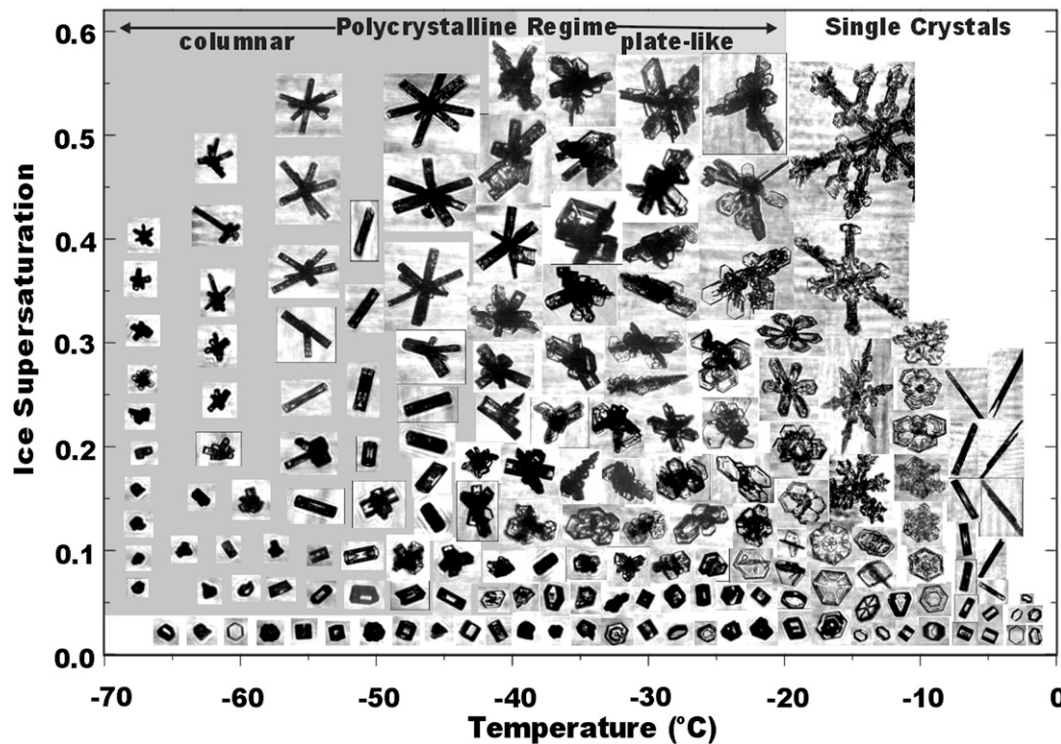
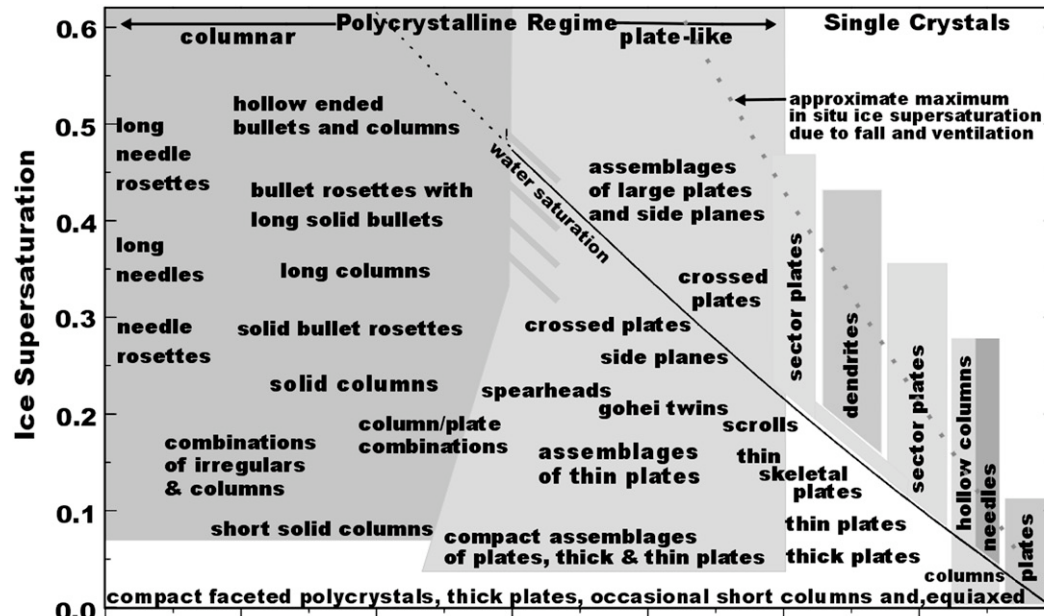
Department of Atmospheric Sciences,
Texas A&M University

Outline

- ◆ Ice Clouds & Science Background
- ◆ Development of Data Library
 - Select representative ice crystal habits
 - Most updated refractive index
 - State-of-the-art computational methods
- ◆ Perspectives & Summary

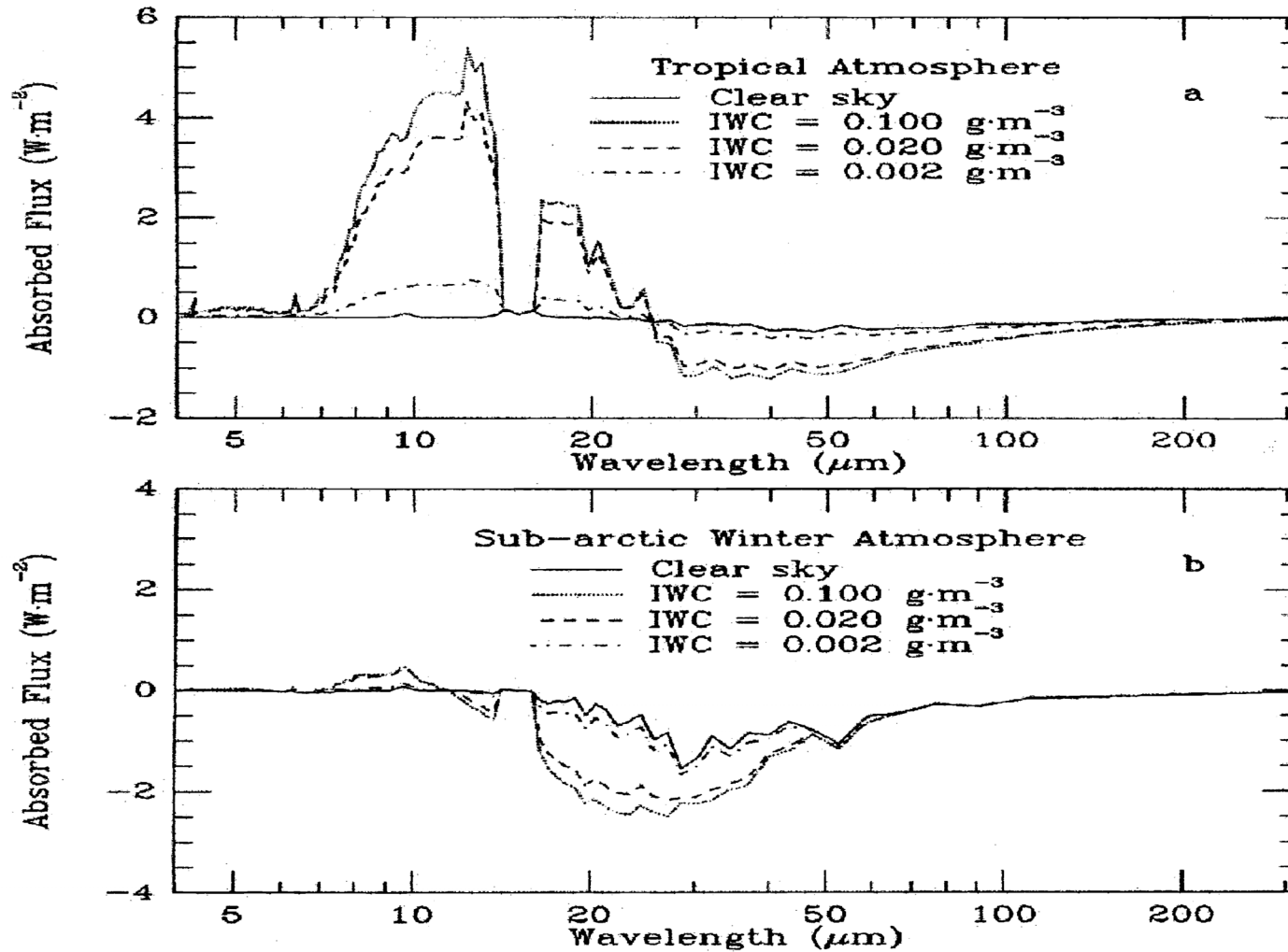
Comprehensive Ice Crystal Habits

Exclusively nonspherical
Endless ice habits



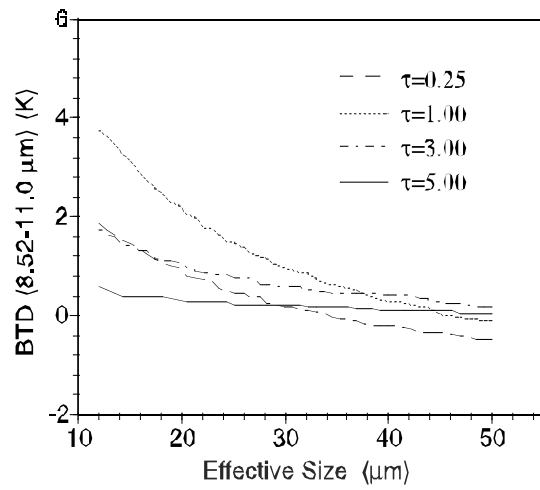
Bailey and Hallett, 2008

Cirrus clouds modulate outgoing far-infrared radiation

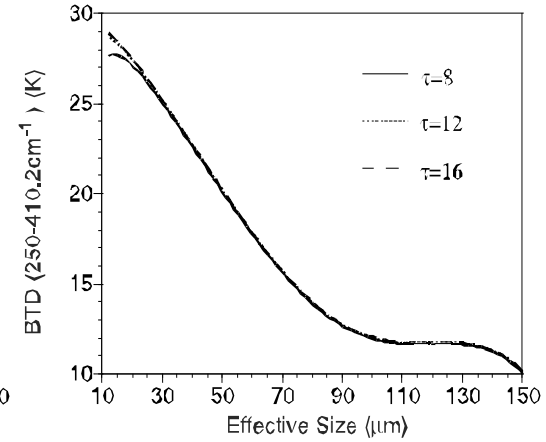
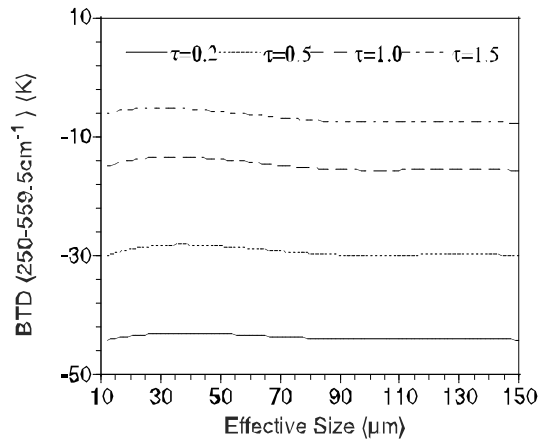


Stackhouse and Stephens, 1991

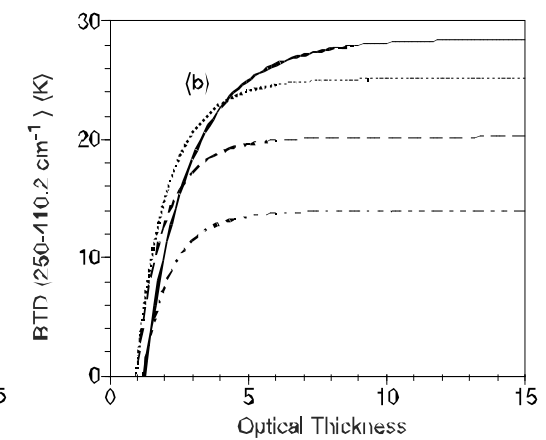
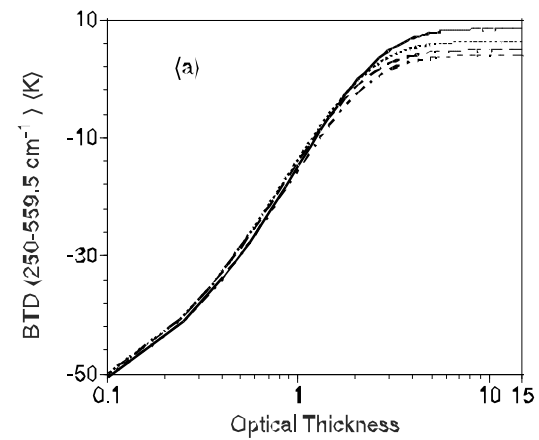
Spectral Signature: Sensitivity



Yang et al (2003)

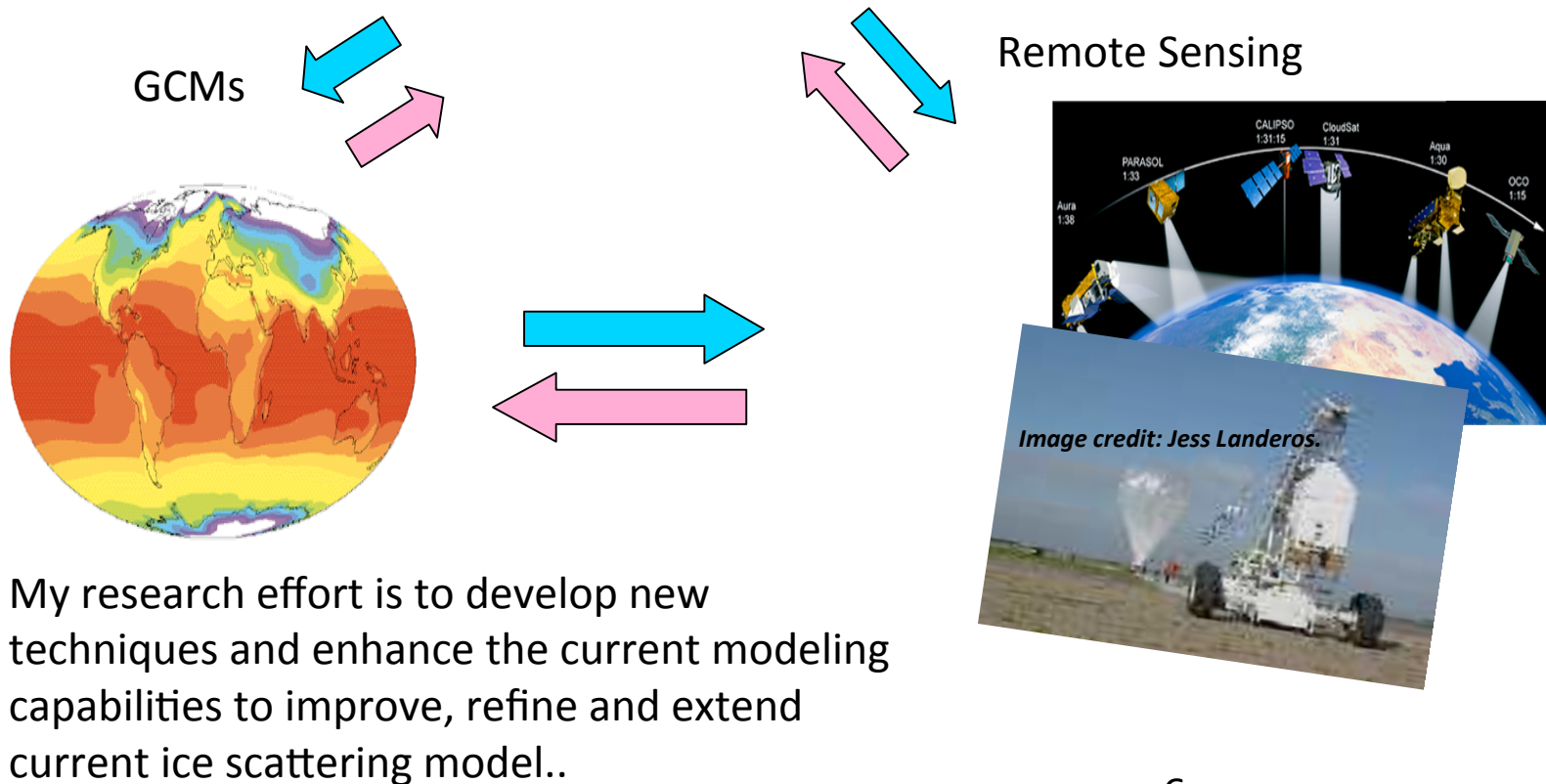


— $D_p=15 \mu\text{m}$ $D_p=30 \mu\text{m}$ - - - $D_p=50 \mu\text{m}$ - - - $D_p=80 \mu\text{m}$

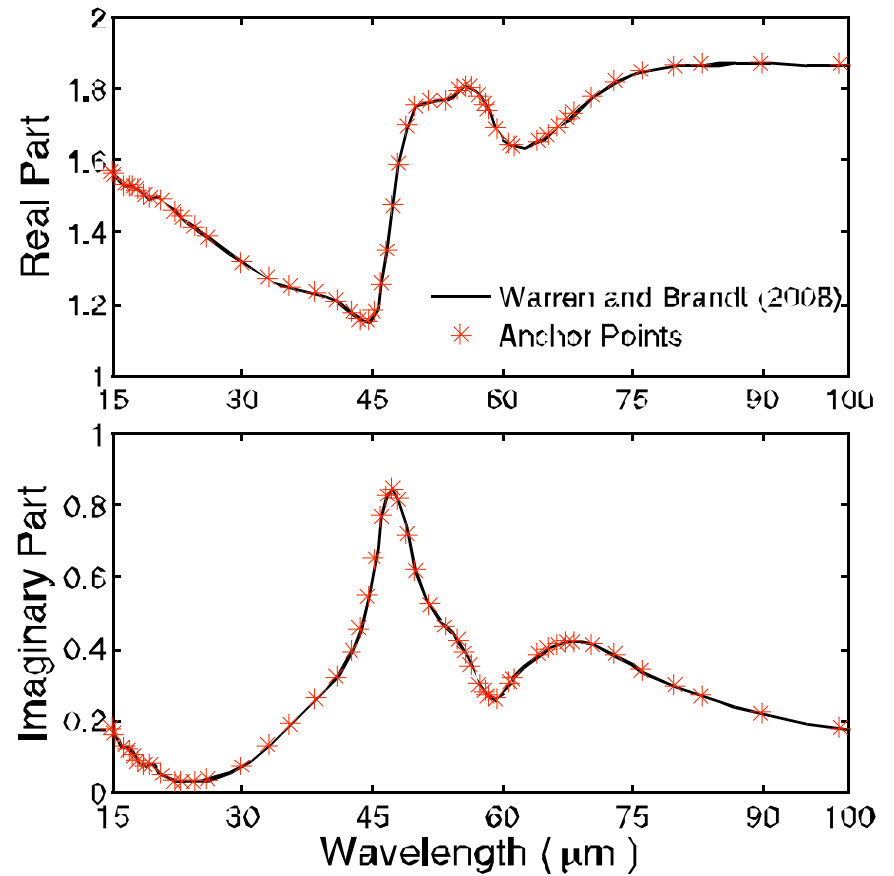
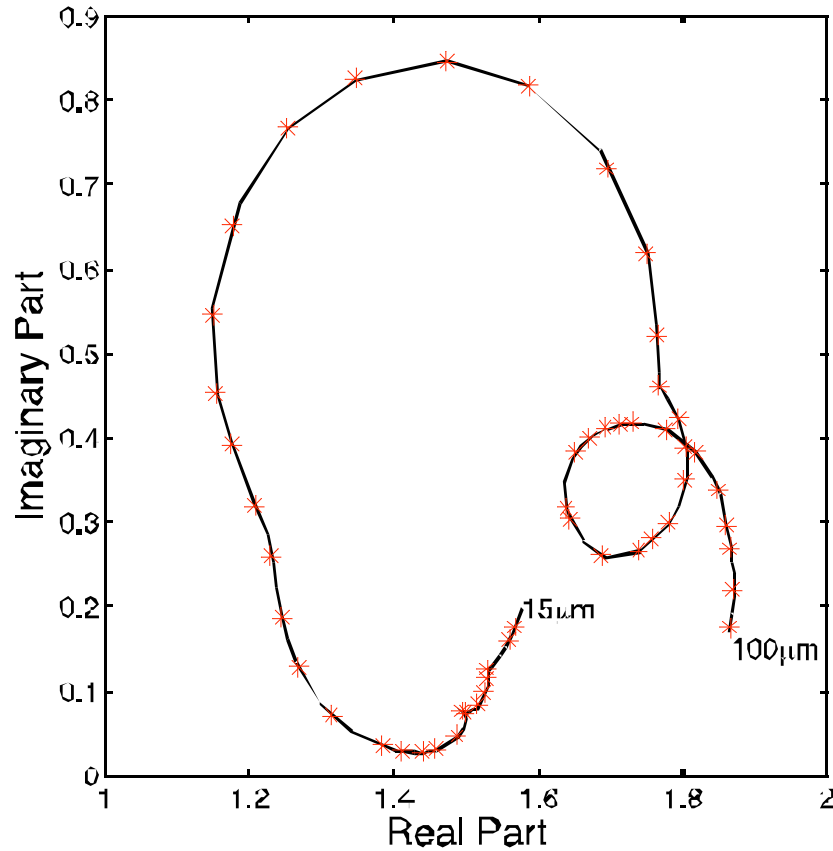


Ice Cloud Model

Ice Cloud Scattering Model

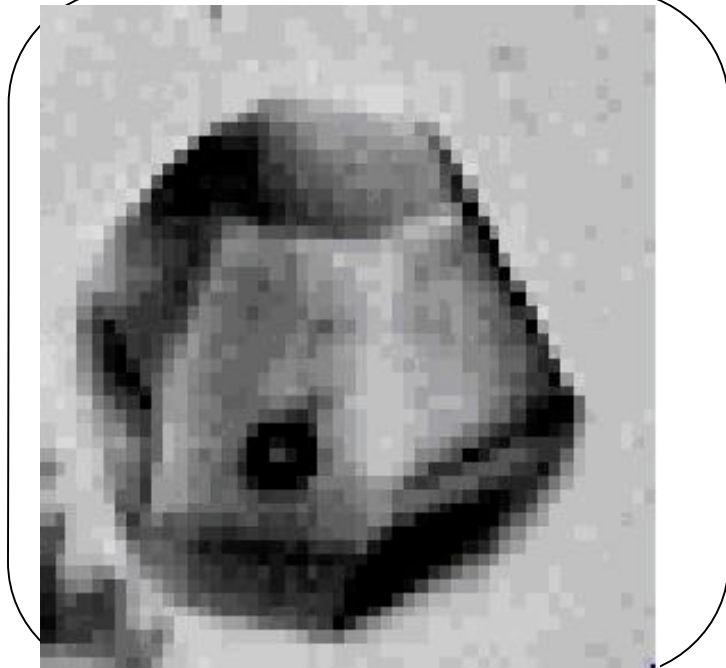
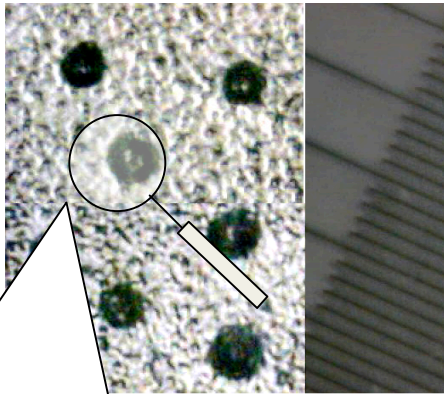


Refractive Index: 15-100 μm

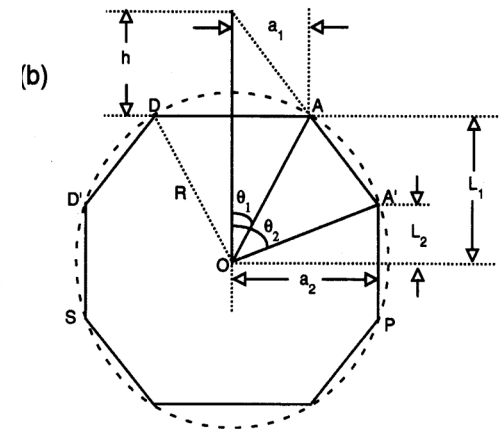
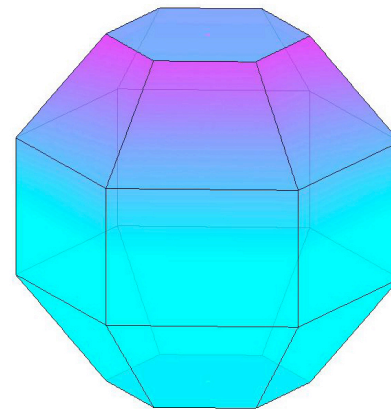


Warren & Brandt (2008)

Ice Crystal Habit: Droxtal



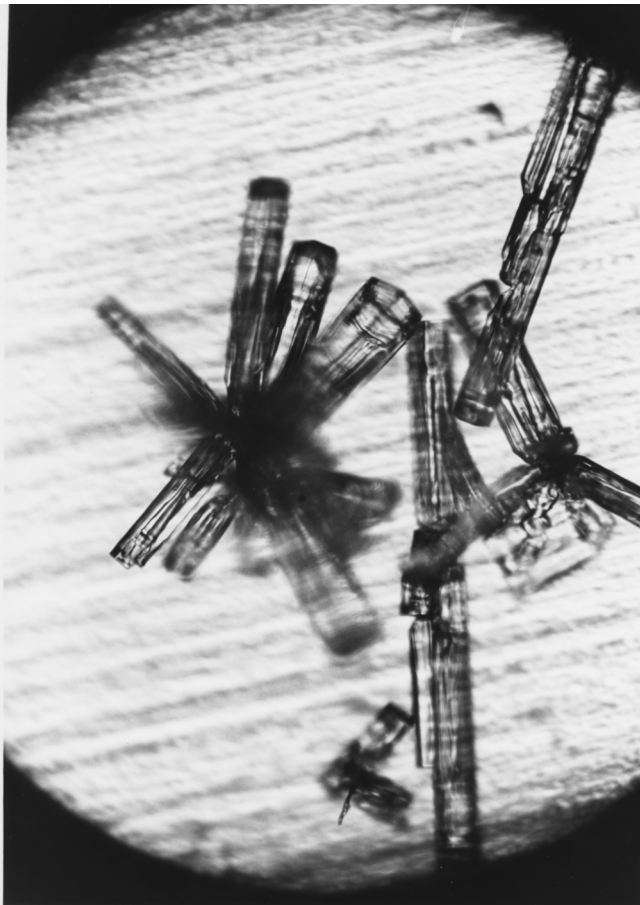
- Dominant in the uppermost portions of cirrus clouds
- 20-Faced Polyhedron
 - 12 isosceles trapezoid
 - 6 rectangular
 - 2 hexagonal



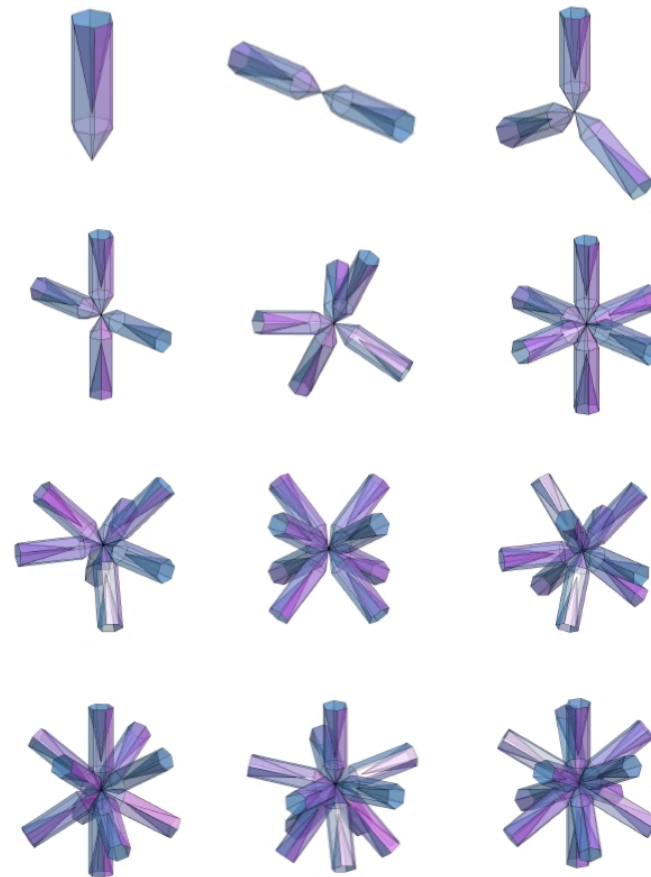
Source : Andy Heymsfield

Ice Crystal Habit: Hollow Bullet Rosettes

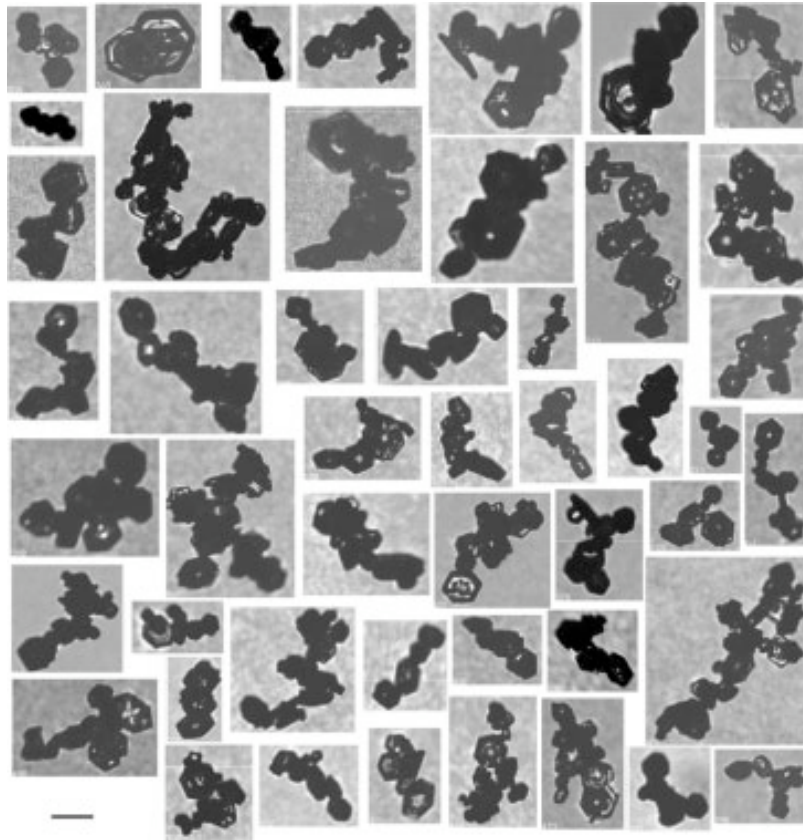
Photo from Stephen G.
Warren



Yang et al (2008)

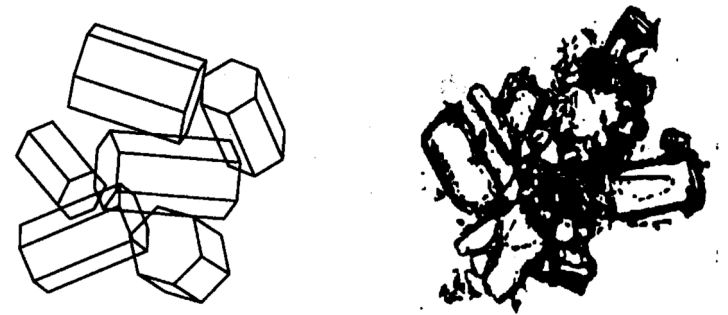
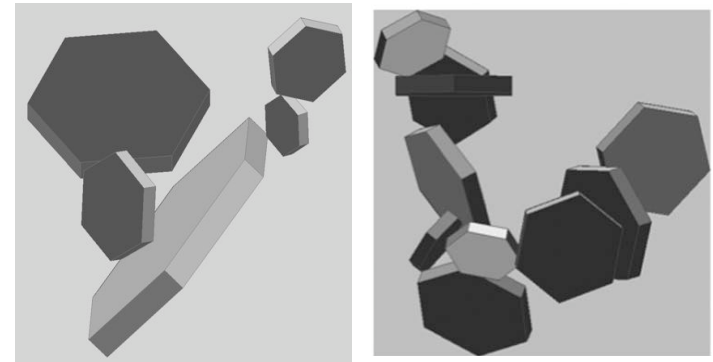


Ice Crystal Habit: Aggregates



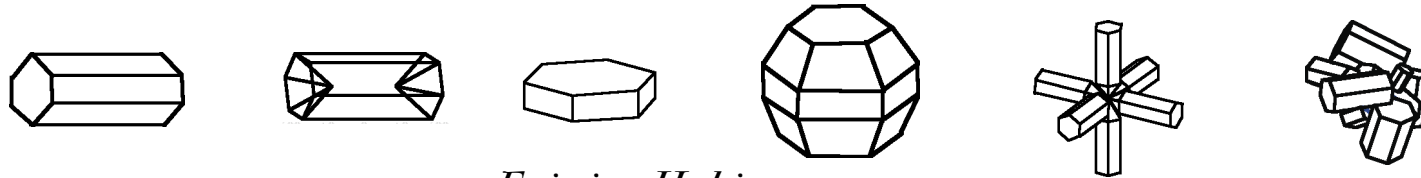
CPI Imag (Um & McFarquhar, 2009)

(Xie, 2011)



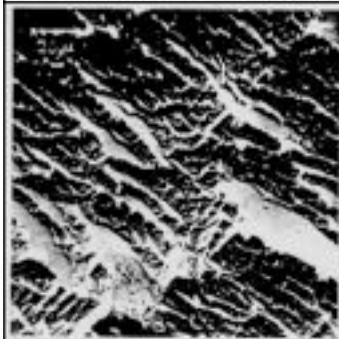
Compact aggregate of columns
(Yang et al 1998)

Ice Crystal Habits



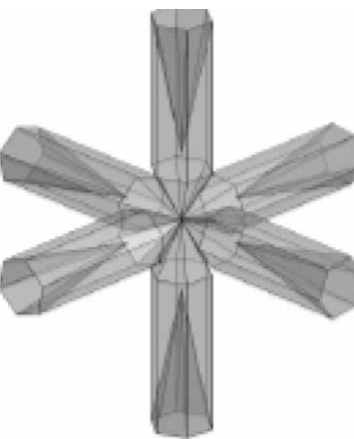
Existing Habits

Incorporated Roughness Condition



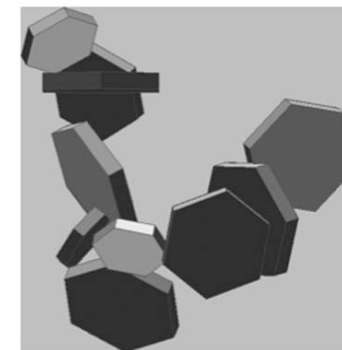
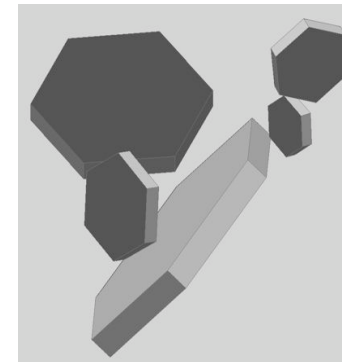
Surface roughness were observed for single crystals and polycrystalline ice by using an electronic microscope. Images adapted from Cross, 1968!

New Habits



Yang et al. (2008)

Xie et al. (2011)



Essence of Radiative Interaction

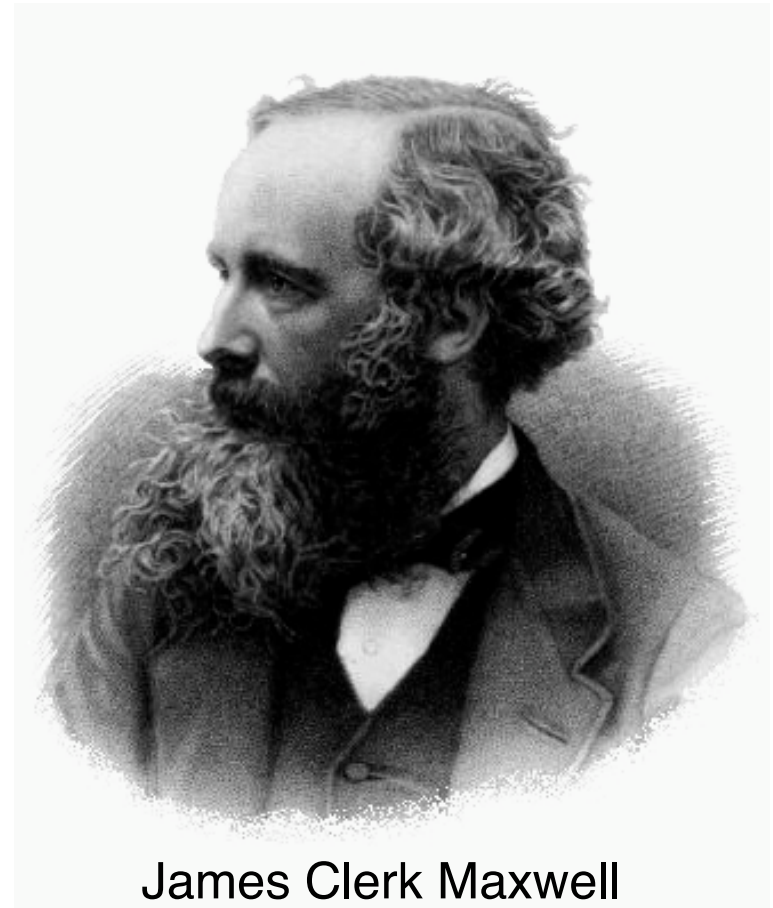
Electromagnetic wave

$$\nabla \times \mathbf{H} = \frac{\varepsilon}{c} \frac{\partial \mathbf{E}}{\partial t}$$

$$\nabla \times \mathbf{E} = -\frac{\mu}{c} \frac{\partial \mathbf{H}}{\partial t}$$

$$\nabla \cdot \mathbf{E} = 0$$

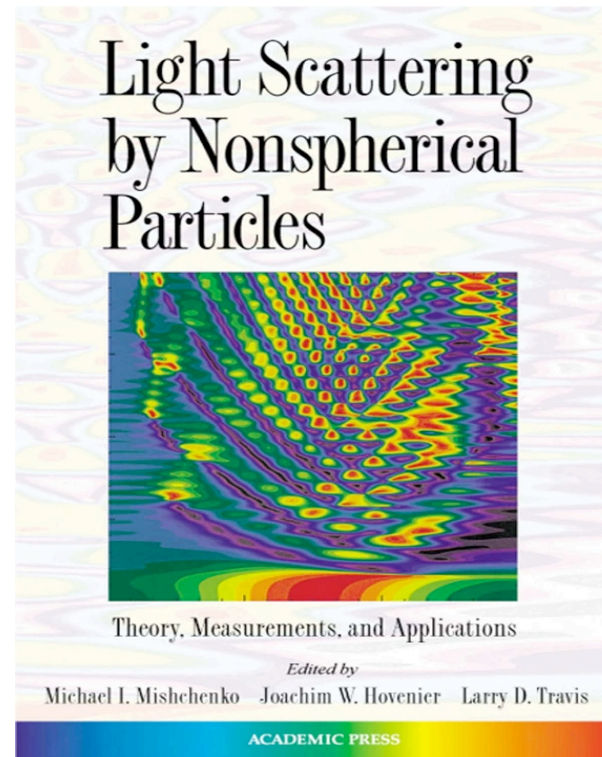
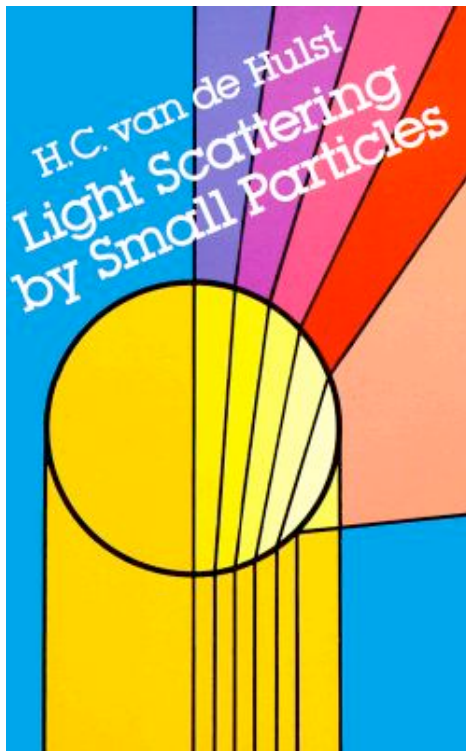
$$\nabla \cdot \mathbf{H} = 0$$



James Clerk Maxwell
(13 June 1831 – 5 November 1879)

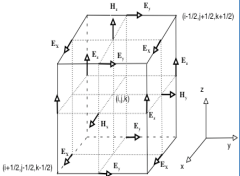
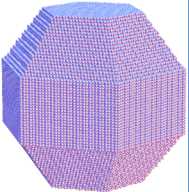
Solution of Maxwell's Equations & Mathematically Equivalents

Van de Hulst (1957)



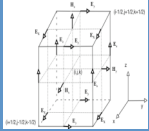
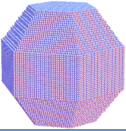
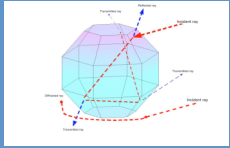
A Combination of Methods

FDTD/DDA

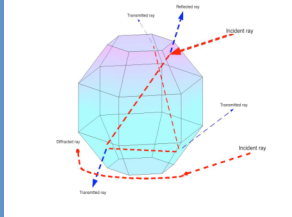



Yee grid (Yee, 1966) Dipoles

FDTD / ADDA + IGOM

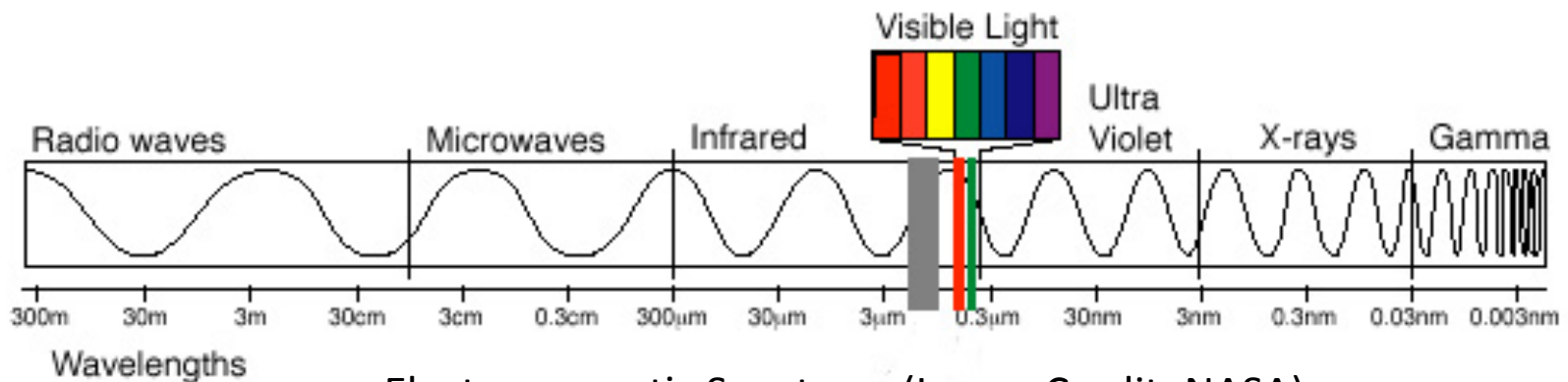
IGOM



Ray tracing



Size Parameter



Electromagnetic Spectrum (Image Credit: NASA)

Discrete-Dipole-Approximation (DDA) Method

SCATTERING AND ABSORPTION OF LIGHT BY NONSPHERICAL DIELECTRIC GRAINS

EDWARD M. PURCELL AND CARLTON R. PENNYPACKER

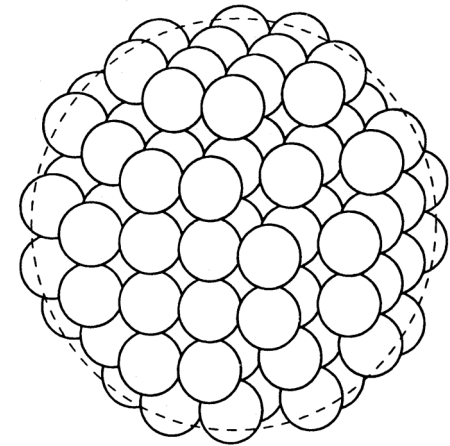
Lyman Laboratory, Harvard University

Received 1973 June 4; revised 1973 July 6

ABSTRACT

A method is described for calculating approximate extinction, absorption, and scattering cross-sections for dielectric grains of arbitrary shape, with dimensions comparable to or smaller than the wavelength of the incident radiation. The grain is modeled by an array of N polarizable elements in vacuum. The elements are located on a simple cubic lattice, and N is of order of magnitude 100. The polarizability of an element is such that an unbounded array would exhibit, according to the Clausius-Mossotti relation, the bulk dielectric constant of the grain material. The complex vector amplitude of each oscillator in the array, which is driven by the field of the incident wave and the fields of all the other oscillators, is determined by an iterative procedure. From the N amplitudes all the cross-sections, including differential cross-sections, are obtained. The method was tested by comparing the cross-sections computed for "spherical" clusters of oscillators with exact Mie theory values for the corresponding dielectric spheres. Computed cross-sections are presented for five different grain shapes and three different complex refractive indices.

Subject headings: interstellar matter — opacities — polarization



ApJ 333 848 (1988)

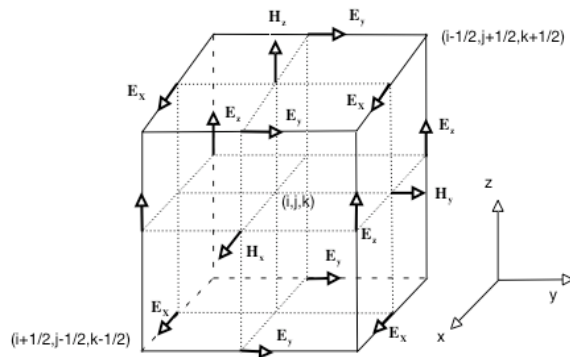
$$\vec{E}(\vec{r}) = \vec{E}_{inc}(\vec{r}) + \int_V d^3r' \vec{G}(\vec{r}, \vec{r}') \chi(\vec{r}') \vec{E}(\vec{r}')$$

$$\vec{G}(\vec{r}, \vec{r}') = \left[\vec{I} + \frac{1}{k^2} \nabla \nabla \right] \frac{\exp(ik|\vec{r} - \vec{r}'|)}{4\pi|\vec{r} - \vec{r}'|}$$

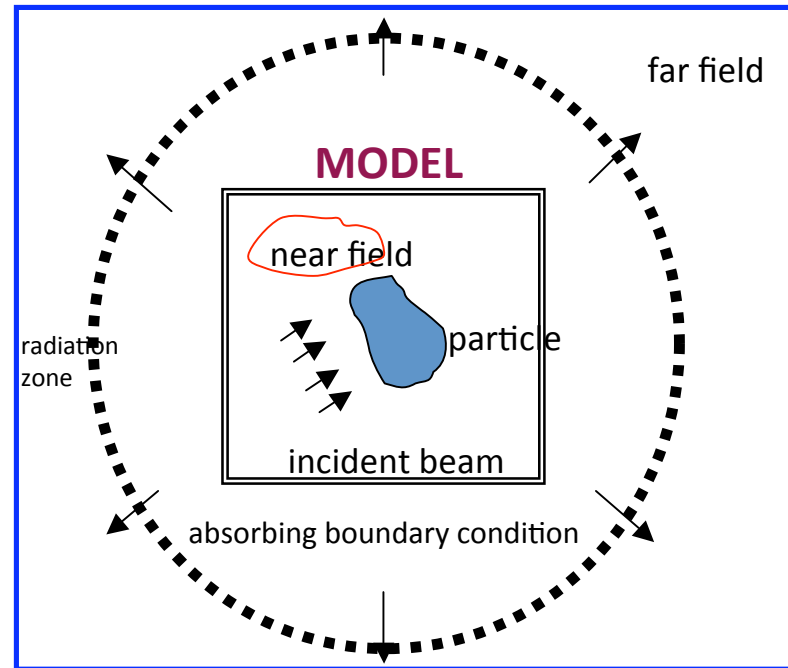
- Efficient parallel-version codes (DDSCAT, ADDA) have been designed (Draine, 2006, ;Yurkin, 2006).

Time-Domain Methods

$$\begin{aligned}\nabla \times \vec{\Pi}(\vec{r}, t) &= \frac{\varepsilon}{c} \frac{\partial \vec{E}}{\partial t}, \\ \nabla \times \vec{E}(\vec{r}, t) &= -\frac{\mu}{c} \frac{\partial \vec{H}}{\partial t}\end{aligned}$$



Yee grid (Yee, 1966)



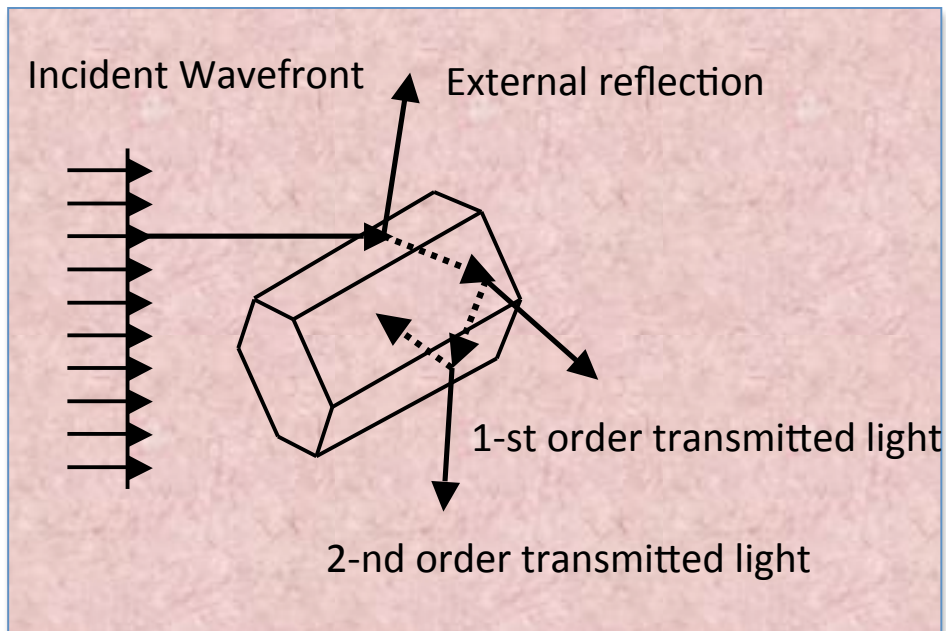
Finite-difference time-domain (FDTD):

$$\frac{\partial E_z^{sca}(r_{i+1/2, j, k}, t_n)}{\partial y} = \frac{1}{\Delta y} \left[E_z^{sca}(r_{i+1/2, j, k}, t_n) - E_z^{sca}(r_{i+1/2, j-1/2, k}, t_n) \right]$$

Pseudo-spectral time-domain (PSTD):

$$\frac{\partial E_z^{sca}(r_{i+1/2, j, k}, t_n)}{\partial y} = F^{-1} \left\{ -ik_y F \left[E_z^{sca}(r, t_n) \right] \right\}$$

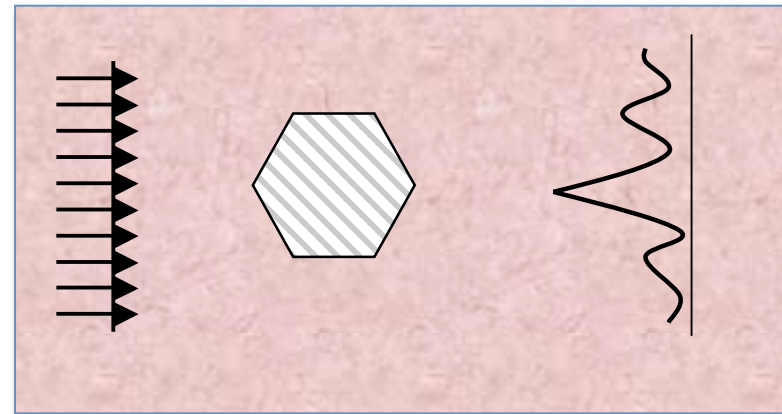
Geometric-Optics Method



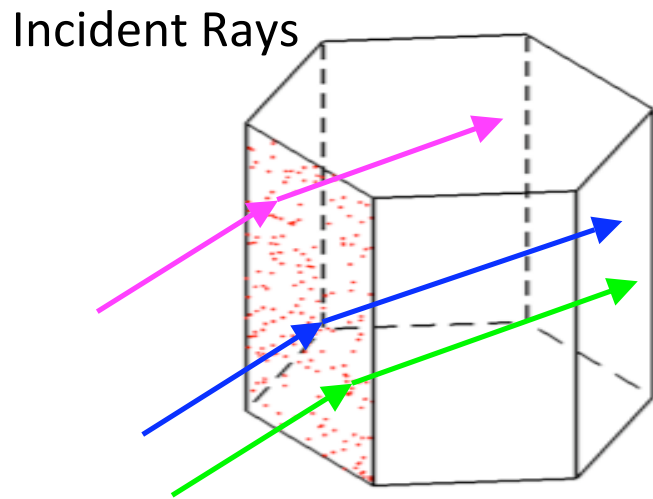
Basic Principle

Ray / Beam / "Photon",
Intensity, Polarization, Phase
Ray/Beam Tracing
Snell's Law, Fresnel Formulas
Fraunhofer Diffraction

In a more rigorous sense, this method is of
physics-geometric optics hybrid (PGOH).

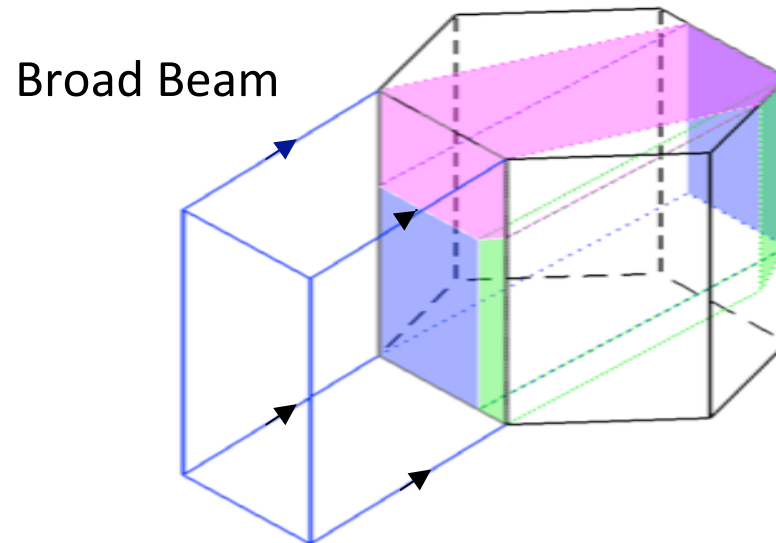


Ray/Beam-tracing Algorithm



Conventional ray-tracing

(a)



Beam-splitting technique

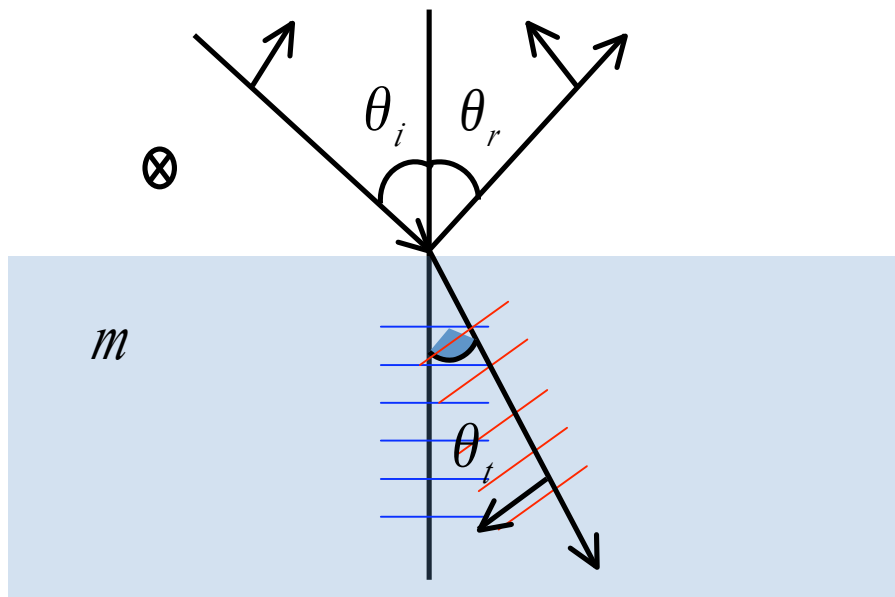
(b)

(a) Larger size parameter, larger number of rays

(b) Number of beams is irrelevant to size parameter.

Similar studies: Popov, 1996; Borovoi, 2003

Inhomogeneous Plane Waves



$$\sin \theta_i = m \sin \theta_t$$

Effective refractive index (Yang and Liou, 2009)

$$N_r = \sqrt{\left\{ m_r^2 - m_i^2 + \sin^2 \theta_i + \sqrt{(m_r^2 - m_i^2 - \sin^2 \theta_i)^2 + 4m^2 m_i^2} \right\} / 2}$$

$$N_i = \cos \theta_t \sqrt{\left\{ -(m_r^2 - m_i^2 - \sin^2 \theta_i) + \sqrt{(m_r^2 - m_i^2 - \sin^2 \theta_i)^2 + 4m^2 m_i^2} \right\} / 2}$$

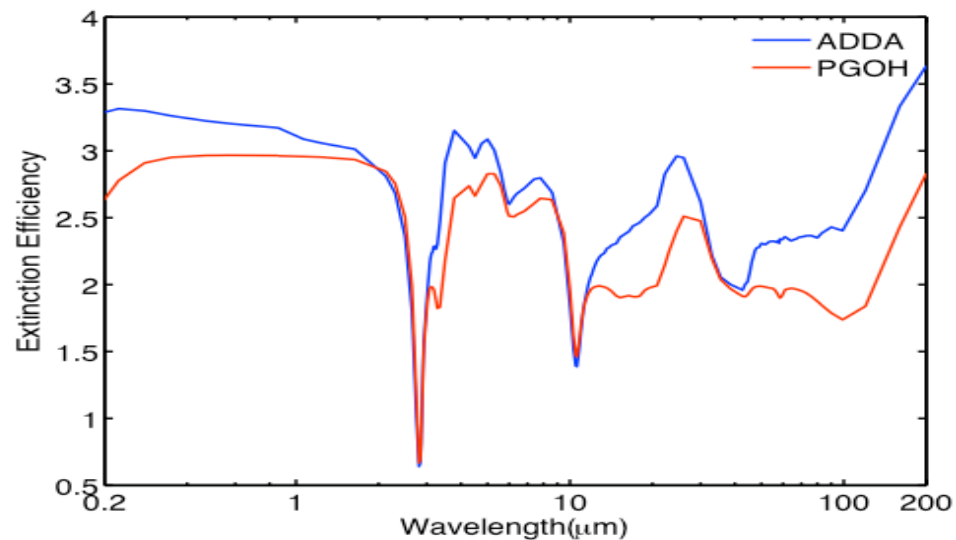
$$E \exp(-kN_i l) \exp(ikN_r \hat{e}_t \cdot r)$$

$$\sin \theta_i = N_r \sin \theta_t$$

Edge-Effect

- Question:

How to understand the difference between extinction efficiency factors simulated from the PGOH and exact methods (e.g. DDA)?



Hexagonal particle. Size parameter is 15.

Localization Principle

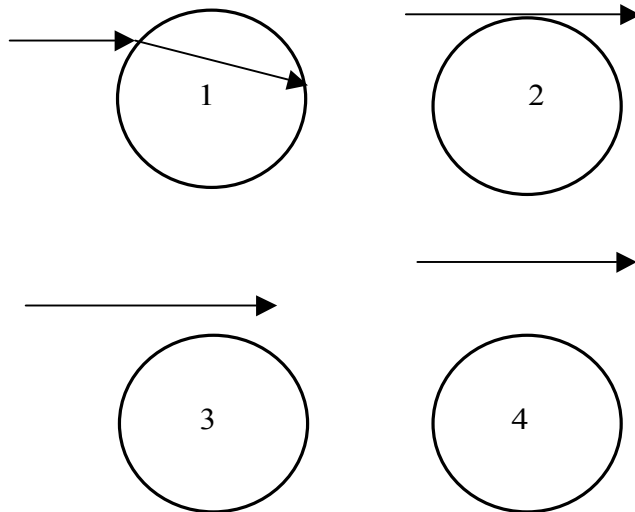
Mie formula:

$$S_1 = \sum_{n=1}^{\infty} \frac{2n+1}{n(n+1)} (a_n \pi_n + b_n \tau_n)$$

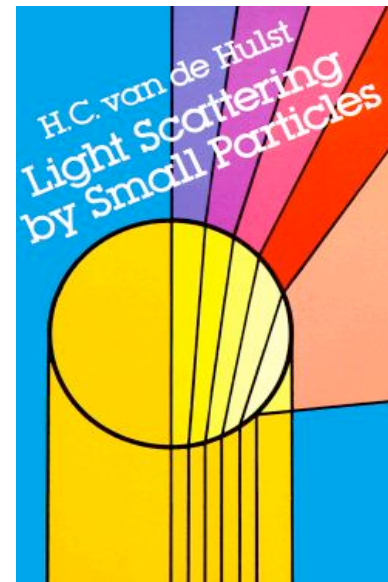
$$S_2 = \sum_{n=1}^{\infty} \frac{2n+1}{n(n+1)} (a_n \tau_n + b_n \pi_n)$$

A term of the order n corresponds to a ray passing the origin at a distance

$$(n + 1/2) \lambda / 2\pi$$

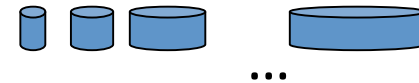
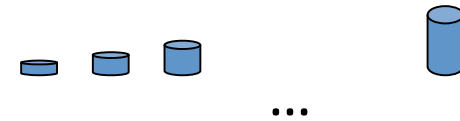
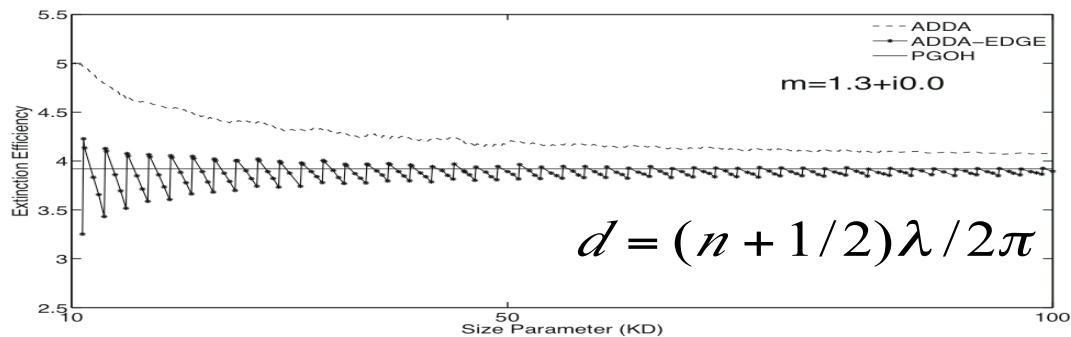
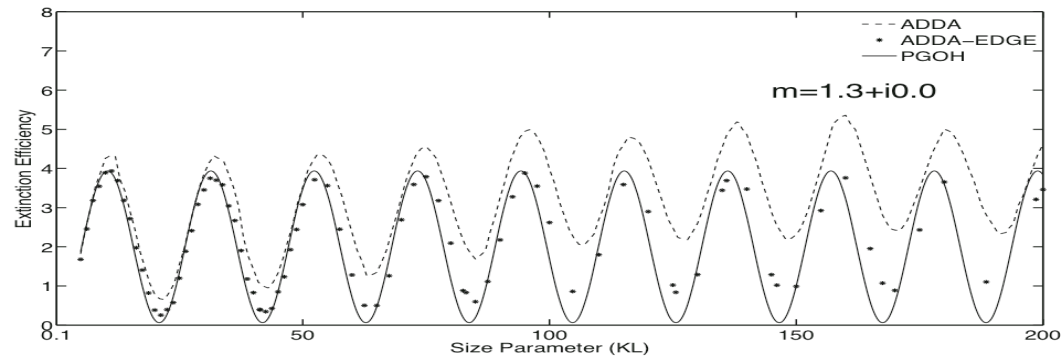
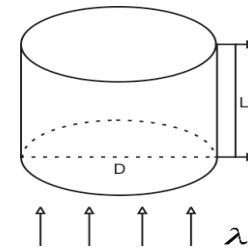


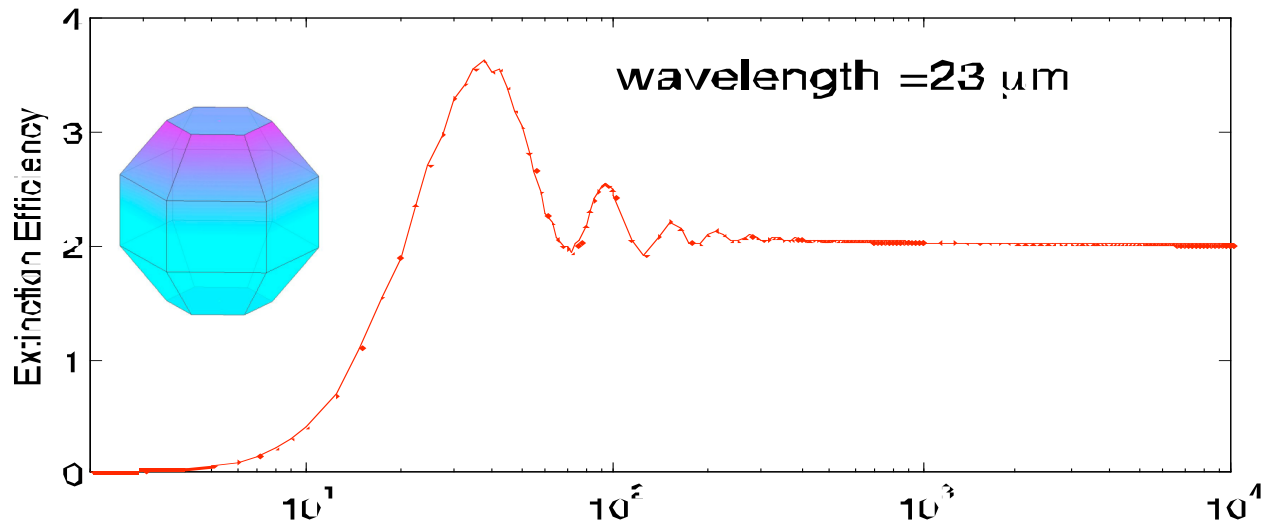
Van de Hulst (1957)



Extinction Efficiency

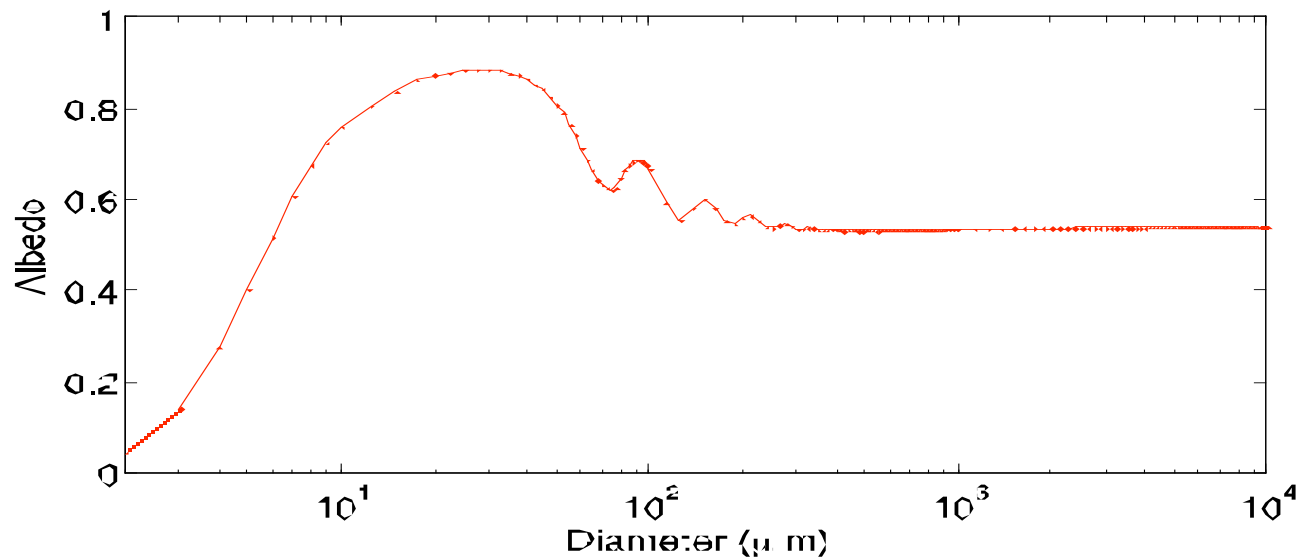
$$Q_{ext} = 2 \operatorname{Re} \left\{ 1 - \frac{4m \exp\{i(m-1)kL\}}{(m+1)^2 - (m-1)^2 \exp(i2mkL)} \right\}$$





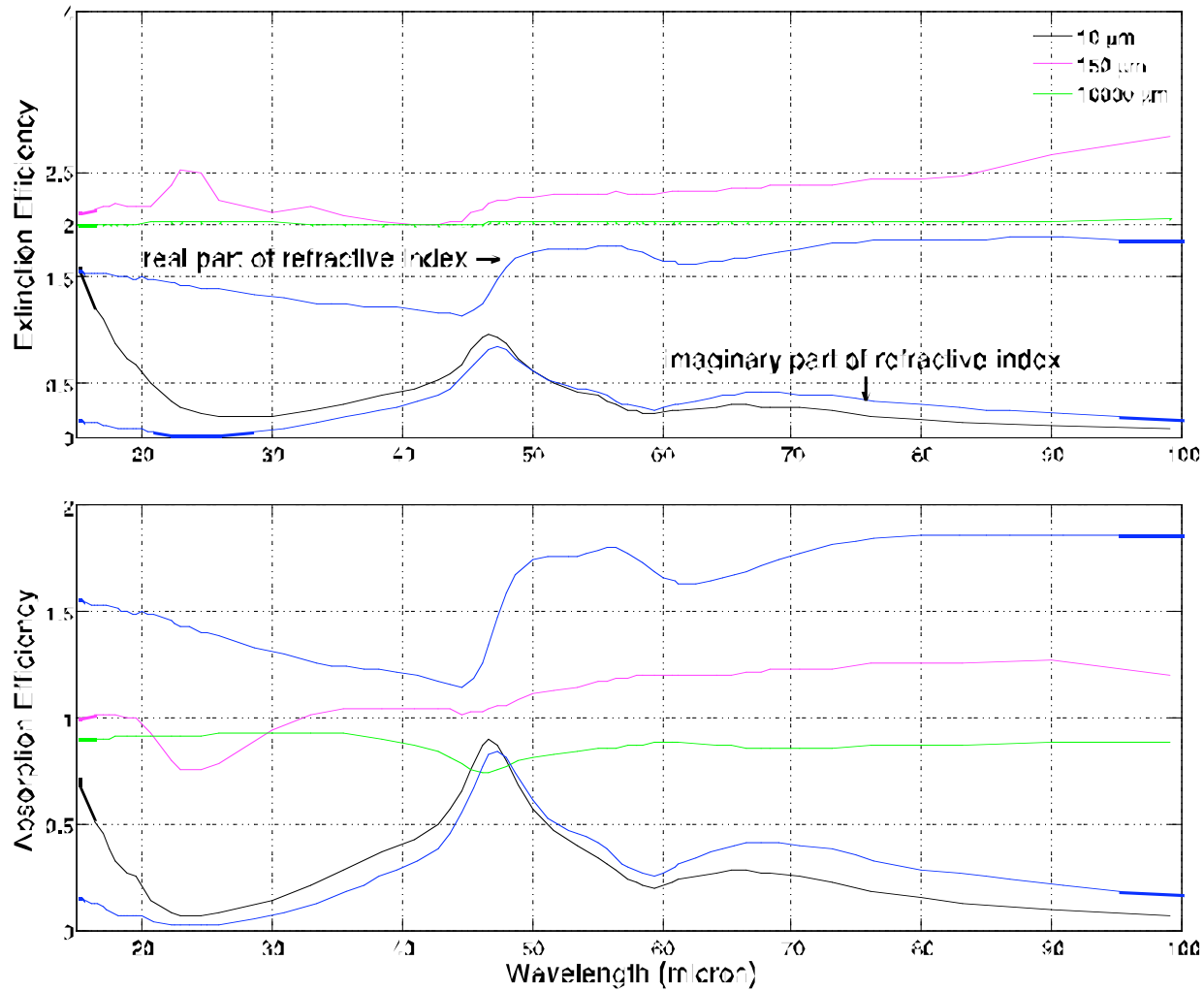
$$Q_{ext, edge} = \frac{f_e}{(kL)^{2/3}}$$

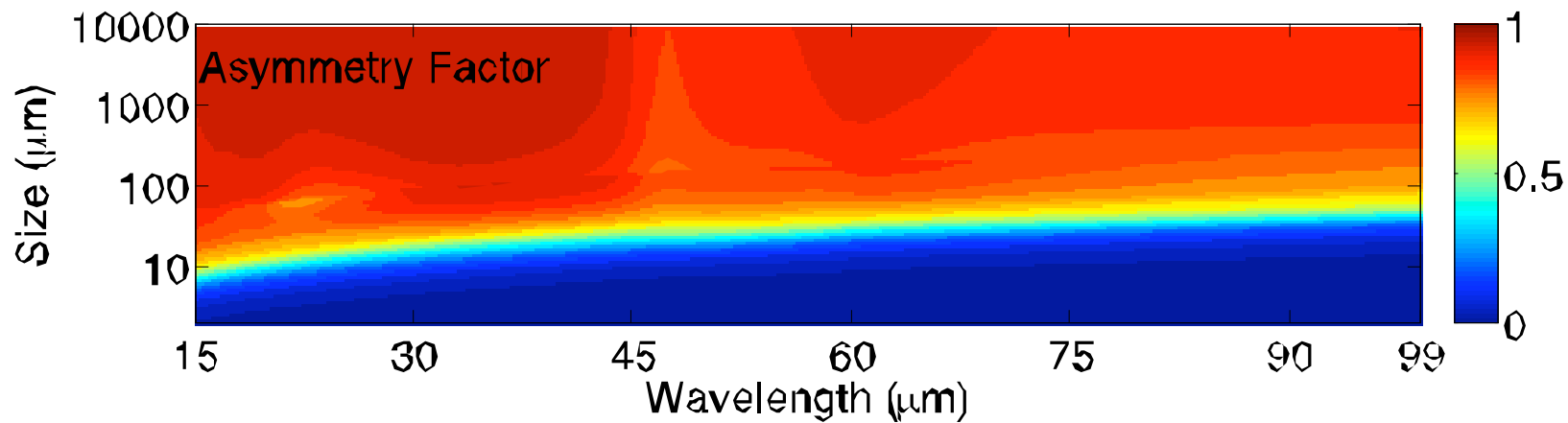
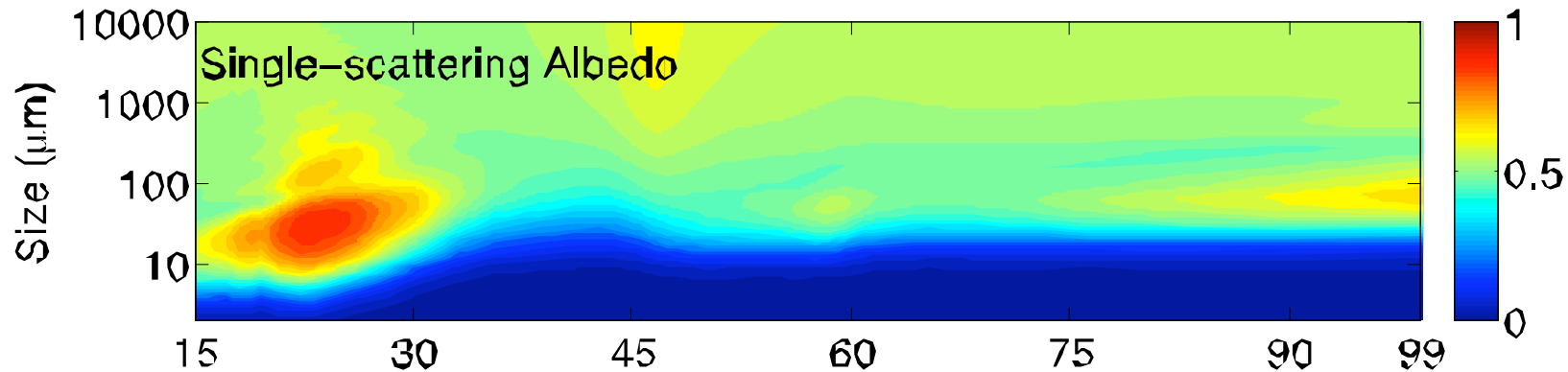
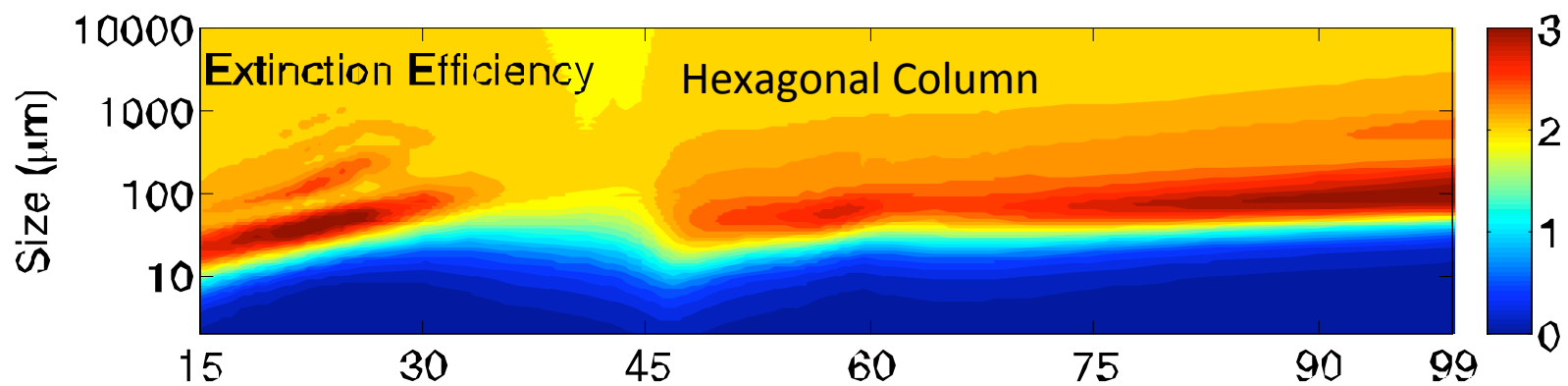
Edge effect incorporated.



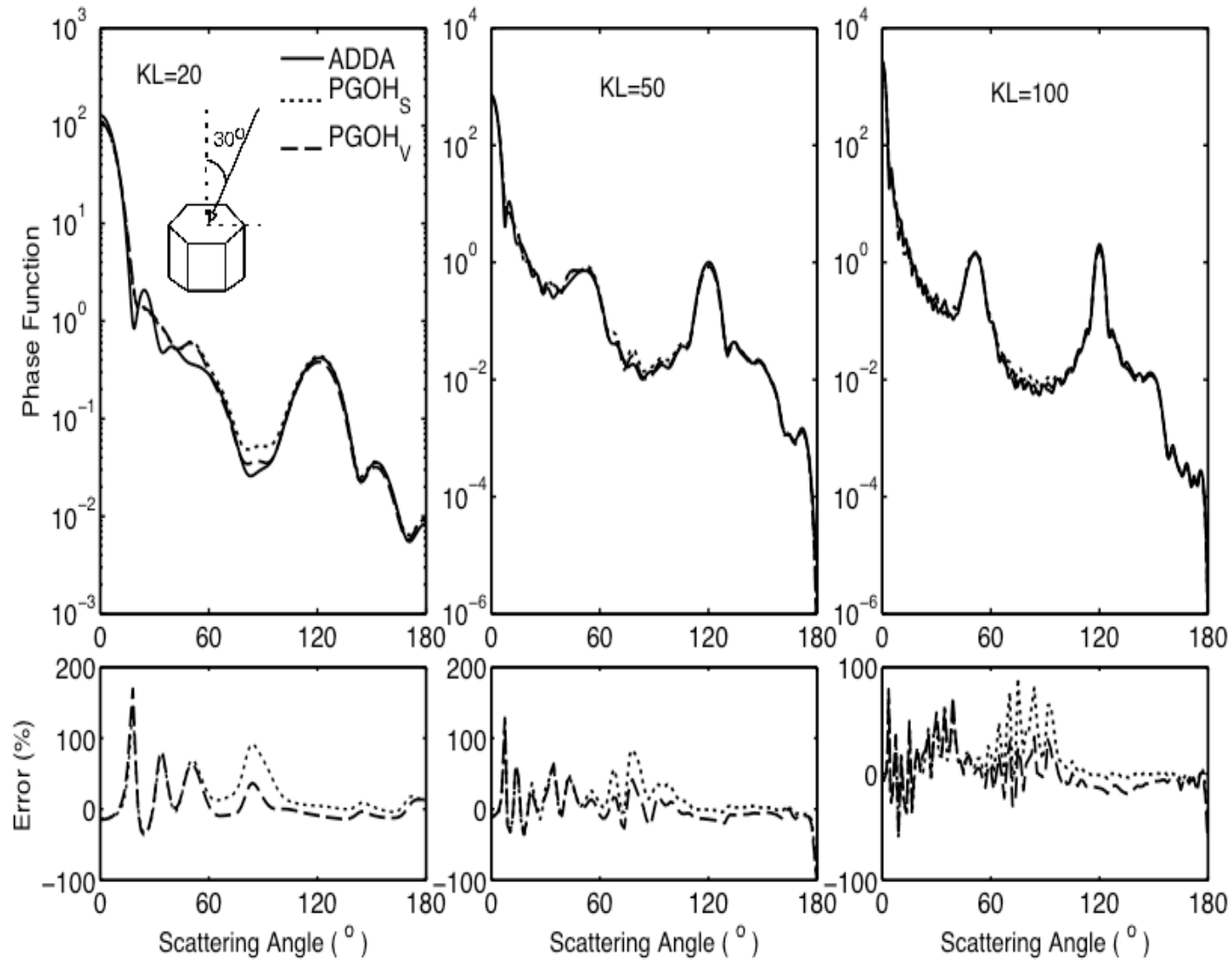
$$Q_{abs, edge} = \frac{f_a}{(kL)^{2/3}}$$

Extinction and Absorption

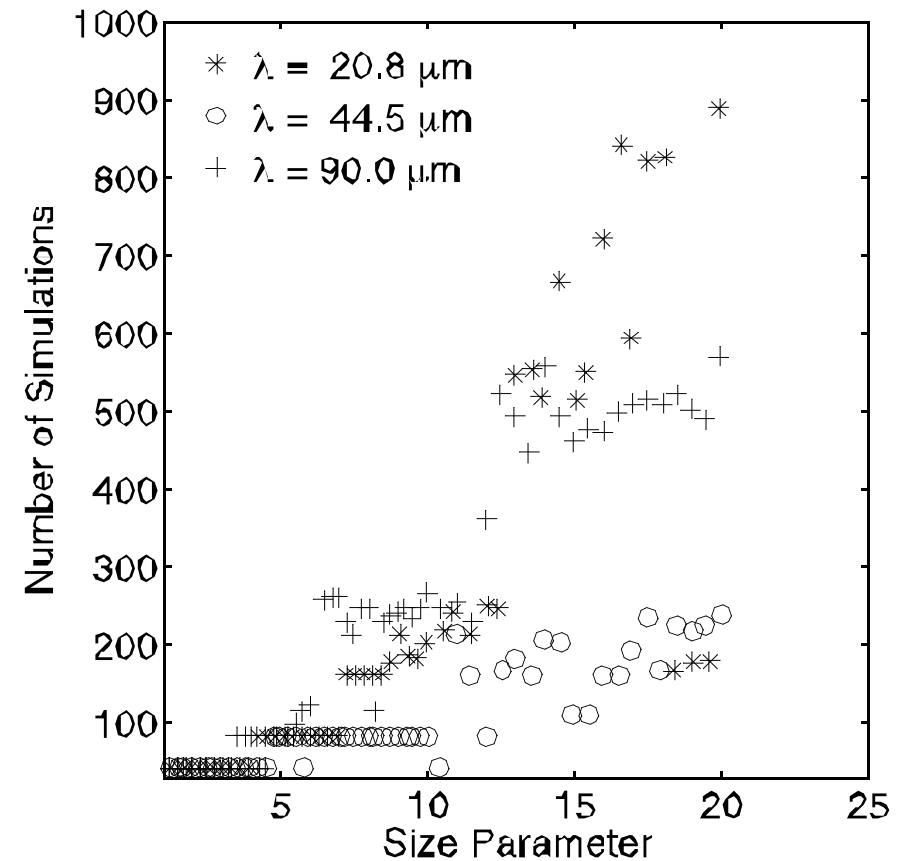
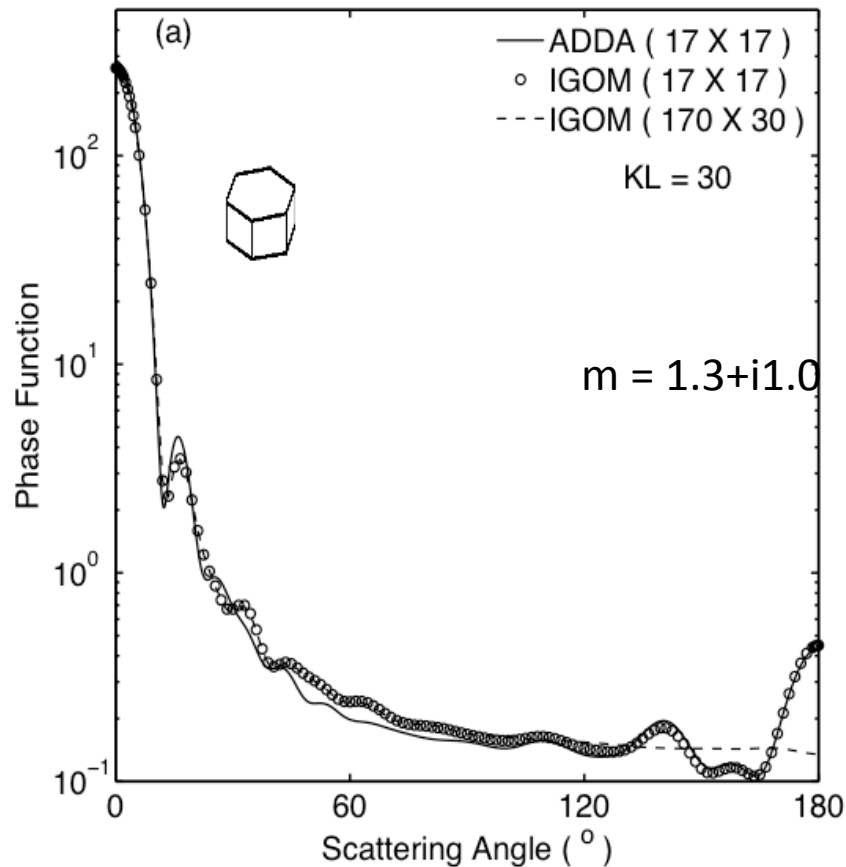




Diffraction and External Reflection

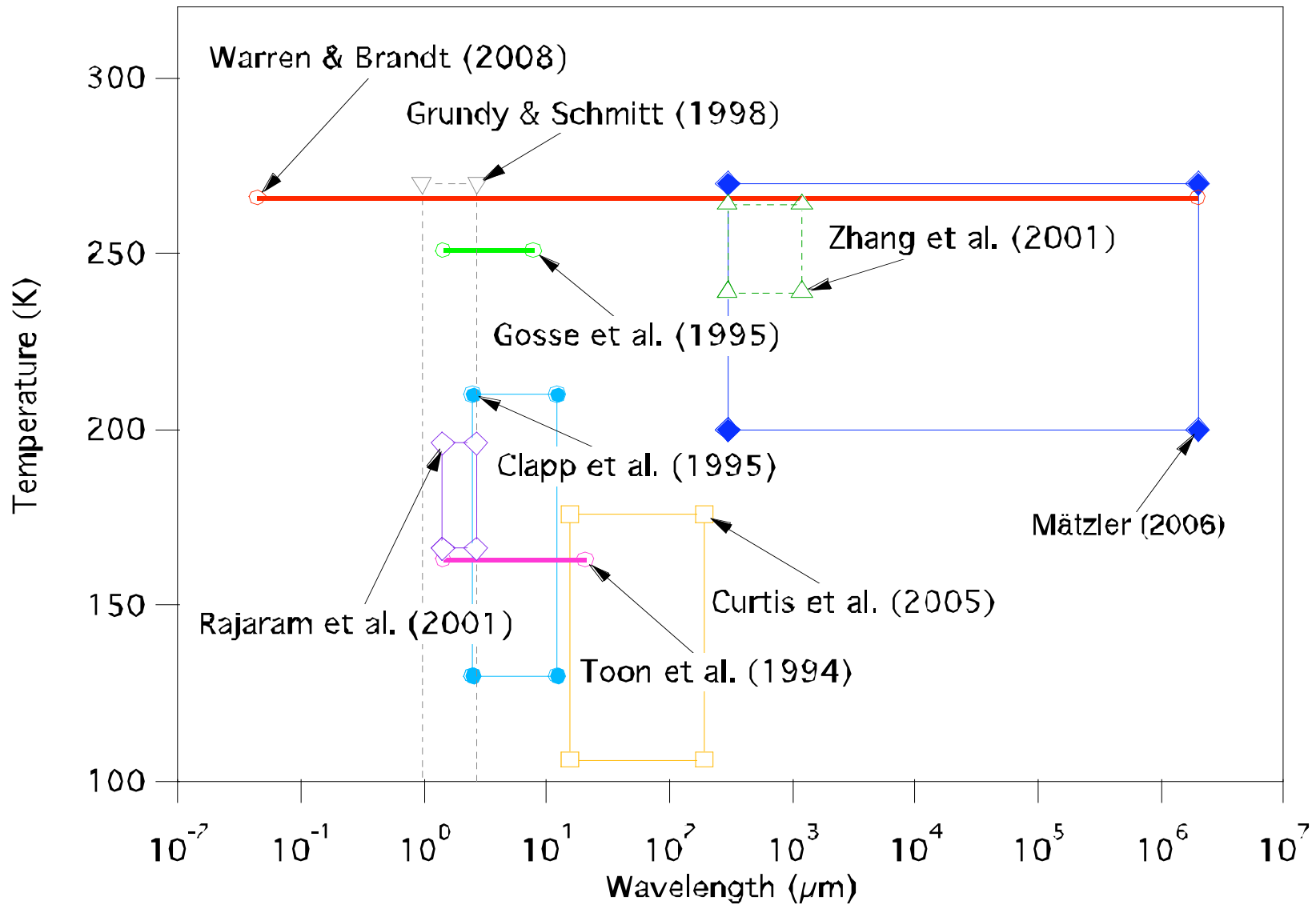


Impact of the Number of particle's Orientations

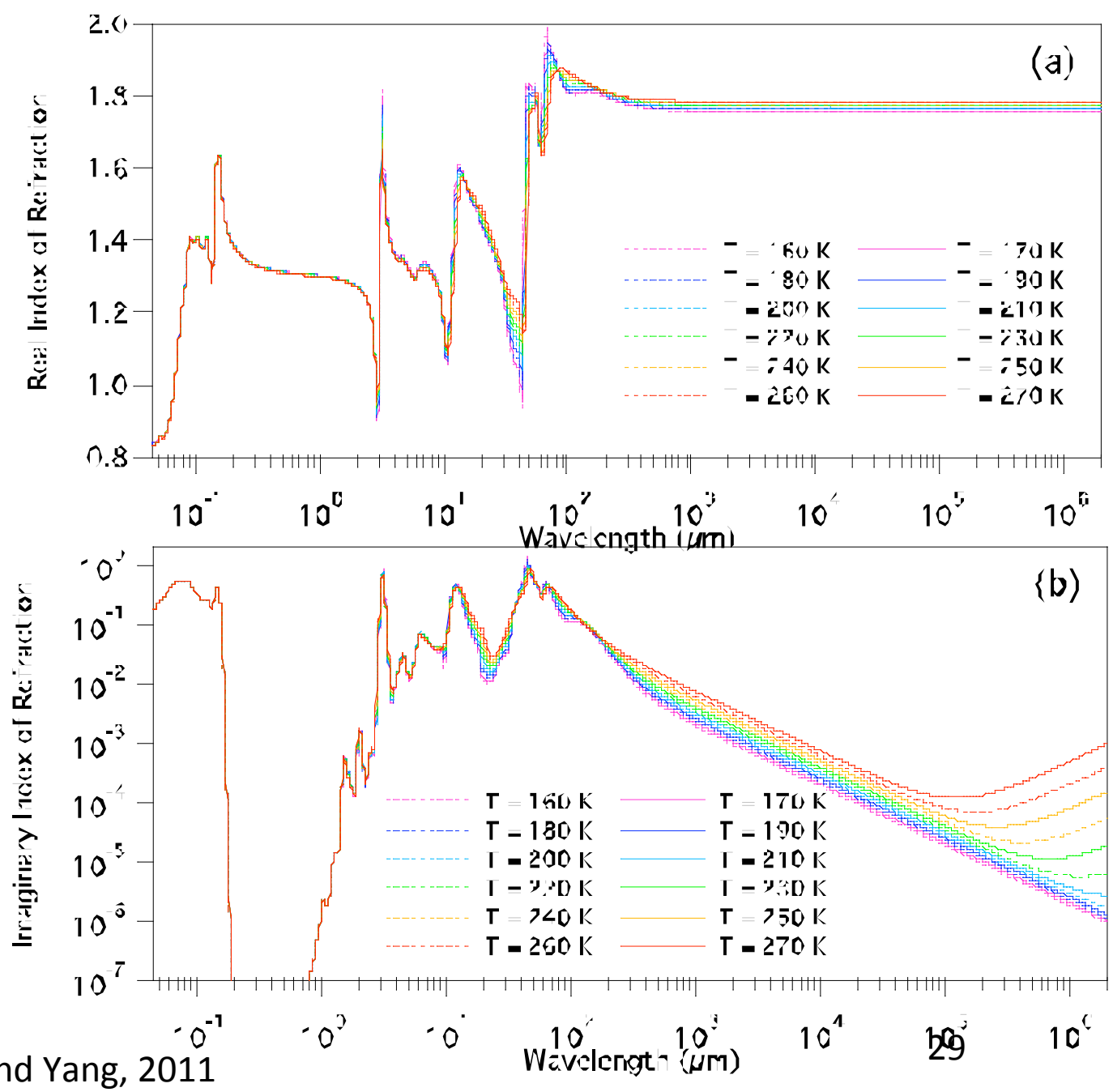


Increasing the number of orientation of model particles!

Temperature dependent refractive index



Iwabuchi and Yang, 2011



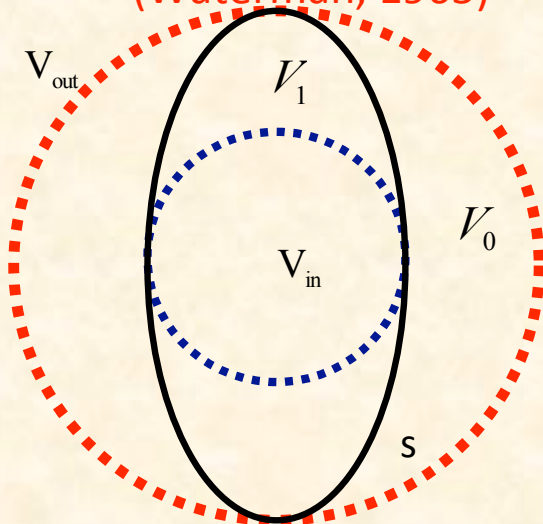
T-matrix

$$E^{inc}(r) = \sum_{n=1}^{\infty} \sum_{m=-n}^n a_{mn} RgM_{mn}(r) + b_{mn} RgN_{mn}(r)$$

$$E^{sca}(r) = \sum_{n=1}^{\infty} \sum_{m=-n}^n p_{mn} M_{mn}(r) + q_{mn} N_{mn}(r)$$

$$\begin{bmatrix} p_{mn} \\ q_{mn} \end{bmatrix} = \sum_{n'=1}^{\infty} \sum_{m'=-n'}^{n'} \begin{bmatrix} T_{mnm'n'}^{11} & T_{mnm'n'}^{12} \\ T_{mnm'n'}^{21} & T_{mnm'n'}^{22} \end{bmatrix} \begin{bmatrix} a_{m'n'} \\ b_{m'n'} \end{bmatrix}$$

Extended Boundary Condition Method
(Waterman, 1965)

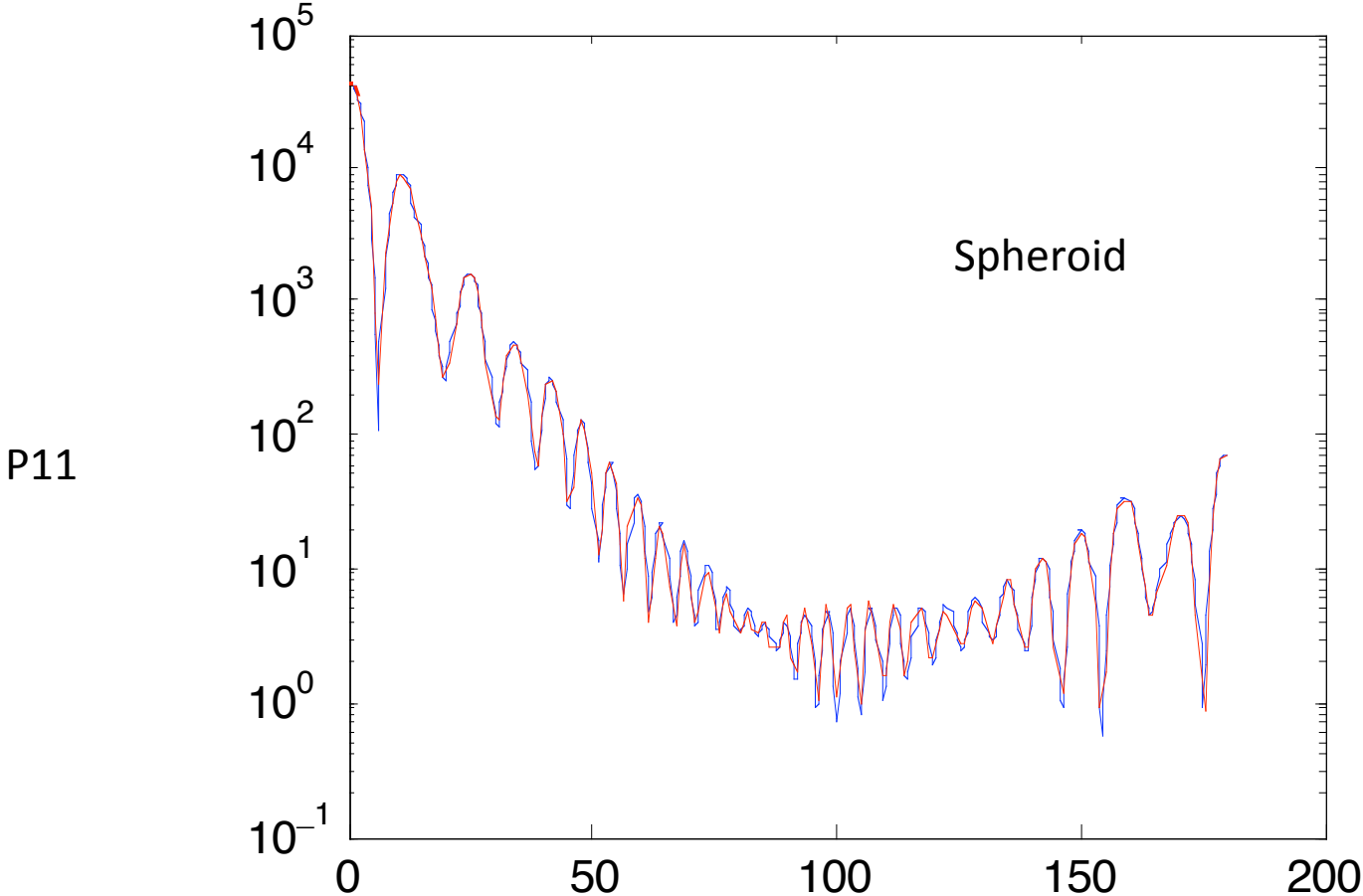


$$E(r) = E^{inc}(r) + \int_V d^3r' G(r, r')(m^2 - 1)E(r')$$

Mie-theory: inscribed sphere
Iteration: remaining volume

Accurate; Stable;
Efficient for random orientation

EBCM T-matrix and new T-matrix



Calculation of EBCM is from Mishchenko's code.

Summary

- Single-scattering property datasets have been computed for the wavelengths from 15~99 μm :
 - Using the updated ice index reported in Warren and Brandt;
 - Improved accuracy of scattering calculations
 - 9 representative ice crystal habits
- Perspectives:
 - Temperature dependence (ice crystal habit and refractive index)
 - Volume-integration T-matrix calculation to further improve the efficiency and accuracy